# GEOMETRY OF MEASURES IN RANDOM SYSTEMS WITH COMPLETE CONNECTIONS 

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#### Abstract

We study new relations between countable iterated function systems, Smale endomorphisms and random systems with complete connections. We prove that the stationary measures of arbitrary countable conformal IFS with place-dependent probabilities, are exact dimensional, and we determine their Hausdorff dimension. Next, we construct a family of fractals in the limit set of a countable IFS with overlaps $\mathcal{S}$, and study the dimension for certain measures supported on these sub-fractals. In particular, we obtain families of measures on sub-fractals, related to the geometry of the system $\mathcal{S}$.


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## 1. Introduction

In this paper we study relations between countable conformal iterated function systems (IFS) with overlaps, Smale endomorphisms, and random systems with complete connections, from the point of view of their geometric and ergodic properties. We provide a surprising common framework for studying measures with certain invariance properties and their dimensions in these systems.

Finite and infinite iterated function systems were studied in many settings, and their invariant measures and limit sets have attracted a lot of interest in the literature, for eg [1], [2], [3], [6], [7], [8], [11], [16], [17], [18], to cite a few. Also there are several connections with the dimension theory for hyperbolic endomorphisms, for eg [20], [24], [12]-[15]. Iterated function systems with place-dependent probabilities (weights) were introduced and studied by Barnsley, Demko, Elton, Geronimo in [1]; see also [2], [3], [23]. Random systems with complete connections were introduced and studied by Iosifescu and Grigorescu in [10] (see also [8]), and are generalizations of the chains with complete connections introduced by Onicescu and Mihoc in [19]. Smale endomorphisms were introduced and studied by the authors in [17].

In the sequel, we first define/recall the notions of countable IFS with overlaps and placedependent probabilities, the notion of random systems with complete connections, and the notion of Smale endomorphisms. We show that a countable IFS with place-dependent probabilities is a particular case of random system with complete connections. Then, in

[^0]Theorem 1.5 we show that, to certain Smale endomorphisms, one can associate random systems with complete connections.

In Section 2, given a countable IFS $\mathcal{S}$ with place-dependent probabilities $\left\{p_{i}(\cdot), i \in I\right\}$ and arbitrary overlaps, we find its stationary measure. Then, in Theorem 2.1 we prove the exact dimensionality of such a stationary measure, and find its Hausdorff (and pointwise, box) dimension. The exact dimensionality largely characterizes the local and global metric properties of the respective measure.

Next, in Section 3, for an arbitrary countable IFS with overlaps $\mathcal{S}$ which satisfies a condition of pointwise non-accumulation, we associate a maximal Smale endomorphism. Using this method, we form a family of random fractals in the limit set $\Lambda$ of $\mathcal{S}$, which correspond to sub-systems of iterations. In Theorem 3.2 we determine the pointwise dimension for a class of invariant measures supported on these random sub-fractals in $\Lambda$. In particular, we obtain families of such measures, which are related to the geometry of the system $\mathcal{S}$.

A finite measure $\mu$ on a metric space $X$ is called exact dimensional (for eg [20], [24]) if there exists a number $\delta \geq 0$ such that for $\mu$-a.e $x \in X$,

$$
\lim _{r \rightarrow 0} \frac{\log \mu(B(x, r))}{\log r}=\delta
$$

In this case it follows that $H D(\mu)=\delta$. Exact dimensionality of a measure $\mu$ is a strong geometric property, and implies that the fractal dimensions of $\mu$ (Hausdorff, pointwise, box dimension) coincide. For finite conformal IFS with overlaps, the exact dimensionality of projections of ergodic measures was proved in [7]. For countable conformal IFS with overlaps, the exact dimensionality of projections of ergodic measures satisfying a finite entropy condition was proved in [16]. The property was studied for hyperbolic diffeomorphisms for eg in [20], [24], and for hyperbolic endomorphisms (non-invertible maps) for eg in [14].

Now, let us recall the three main notions that will be used in the sequel:
First, there is the notion of finite iterated function systems with place-dependent probabilities, introduced in [1]; see also the papers [2], [3], [23]. These iterated function systems are particular cases of the chains with complete connections introduced in [19]; see also [5], [9].
A chain with complete connections is a sequence of random variables $\xi_{1}, \xi_{2}, \ldots$ taking real values in a countable set $\Omega$, where the probability that at step $n, \xi_{n}$ takes value $\omega \in \Omega$, depends on the values taken by all previous random variables. Thus,

$$
P\left(\xi_{n+1}=\omega \mid \xi_{n}=\omega_{0}, \xi_{n-1}=\omega_{-1}, \ldots\right)=P(c, \omega)
$$

where $\omega \in \Omega$ and $c=\left(\ldots, \omega_{-1}, \omega_{0}\right)$ is a trajectory in $\Sigma_{\Omega}^{-}:=\prod_{j \in \mathbb{Z}, j \leq 0} \Omega$. One assumes in general that $P^{1}(c, \omega)=P(c, \omega)$ and for every $n \geq 1$,

$$
P^{n+1}(c, \omega)=\sum_{\omega^{\prime} \in \Omega} P\left(c, \omega^{\prime}\right) P^{n}\left(a\left(c, \omega^{\prime}\right), \omega\right)
$$

where $a\left(c, \omega^{\prime}\right):=\left(\ldots, \omega_{-1}, \omega_{0}, \omega^{\prime}\right)$, for $c=\left(\ldots, \omega_{-1}, \omega_{0}\right) \in \Sigma_{\Omega}^{-}$.

In the sequel, we extend the notion of finite IFS with place-dependent probabilities to countable iterated function systems with overlaps and place-dependent probabilities. Such a system $\mathcal{S}$ is given by continuous functions $\phi_{i}: V \longrightarrow V, \quad i \in I$, defined on a compact set $V \subset \mathbb{R}^{D}$ and indexed by a countable set $I$, and continuous probability functions (weights), $p_{i}: V \longrightarrow[0,1], i \in I$, satisfying

$$
\sum_{i \in I} p_{i}(x)=1
$$

By IFS with overlaps we mean that the sets $\phi_{i}(V), i \in I$ may intersect in any way, thus no Open Set Condition is assumed (see for eg [6], [7], [16]). Assume also that there exists a number $s \in(0,1)$ such that,

$$
\begin{equation*}
\left|\phi_{i}(x)-\phi_{i}(y)\right| \leq s|x-y|, \forall i \in I, x, y \in V \tag{1.1}
\end{equation*}
$$

The countable IFS case is different from the finite case, since the fractal limit set may be non-compact, and many methods from the finite IFS case do not work.

For $x \in V$ and a Borel set $B \subset V$, the probability of transfer from $x$ to $B$ is equal to $P(x, B)=\sum_{i=1}^{\infty} p_{i}(x) \delta_{\phi_{i}(x)}(B)$. We have then the associated transfer operator:

$$
\mathcal{L} g(x)=\int_{V} g(y) P(x, d y)=\sum_{i=1}^{m} p_{i}(x) g\left(\phi_{i}(x)\right)
$$

where $g: V \rightarrow \mathbb{R}$ is measurable. If $M(V)$ denotes the space of finite signed Borel measures on $V$, then the operator $\mathcal{L}^{*}$ adjoint to $\mathcal{L}$, restricted to the space $M(V)$, is given by

$$
\begin{equation*}
\mathcal{L}^{*} \mu(B)=\int P(x, B) d \mu(x)=\sum_{i=1}^{\infty} \int_{\phi_{i}^{-1} B} p_{i}(x) d \mu(x) \tag{1.2}
\end{equation*}
$$

A Borel probability measure $\mu$ on $V$ is called stationary for the above system if

$$
\begin{equation*}
\mathcal{L}^{*} \mu=\mu, \tag{1.3}
\end{equation*}
$$

and attractive if for all probabilities $\nu$ on $V$ and all bounded measurable $g: V \rightarrow \mathbb{R}$,

$$
\lim _{n \rightarrow \infty} \int_{X} g d\left(\mathcal{L}^{* n} \nu\right)=\int_{V} g d \mu
$$

One of the central problems in the theory of chains with complete connections and of IFS with place-dependent probabilities is to find stationary measures and to study their ergodic and metric properties. In many of these results in the finite case, the probability functions $p_{i}, i=1, \ldots, m$, satisfy a Hölder type condition ([1], [2], [10]).

The second main notion is that of random systems with complete connections, which are generalizations of chains with complete connections (for eg [8], [10]).
Definition 1.1 ([10]). A random system with complete connections (or RSCC) is a quadruple $((W, \mathcal{W}),(X, \mathcal{X}), u, P)$ where:
i) $(W, \mathcal{W})$ and $(X, \mathcal{X})$ are measurable spaces,
ii) $u: W \times X \rightarrow W$ is a measurable map, with the product $\sigma$-algebra $\mathcal{W} \times \mathcal{X}$ on $W \times X$,
iii) $P$ is a transition probability function from $(W, \mathcal{W})$ to $(X, \mathcal{X})$, i.e. $P(w, \cdot)$ is a probability on $\mathcal{X}$ for any $w \in W$ and $P(\cdot, A)$ is a random variable on $Z$ for any set $A \in \mathcal{X}$.

We call $W$ the state space and $X$ the index space. $W$ is assumed to be a locally compact and $\sigma$-compact metric space. For $\mathcal{W}$ we take the $\sigma$-algebra generated by open sets of $W$.

An example of RSCC is an urn scheme with replacement. Consider an initial urn $U_{0}$, which contains $a_{j}=a_{j}^{(0)}$ balls of color $j, 1 \leq j \leq m$. If on trial $n \geq 1$ we extract a ball of color $j$, then this ball is replaced together with $d_{j}$ balls of same color (the rest of the balls being left unchanged), hence $a_{i}^{(n)}=a_{i}^{(n-1)}+\delta_{i j} d_{i}, 1 \leq i \leq m$, with $d_{1}, \ldots, d_{m}$ being non-negative integers. Thus the probability of choosing a certain color at step $n$ depends on all previous steps.
Remark 1.2. If in Definition 1.1 the index space $X$ is countable and $\mathcal{X}$ is the algebra of subsets of $X$, and we define the maps $\phi_{i}(w):=u(w, i)$, then we obtain a countable IFS on $W$. Denote $x^{(n)}:=\left(x_{1}, \ldots, x_{n}\right) \in X^{n}, n \geq 1$. By induction define $u^{(n)}: W \times X^{n} \rightarrow W$,

$$
u^{(n+1)}\left(w, x^{(n+1)}\right)=u\left(w, x_{1}\right), n=0, \text { and } u^{(n+1)}\left(w, x^{(n+1)}\right)=u\left(u^{(n)}\left(w, x^{(n)}\right), x_{n+1}\right), n \geq 1
$$

We can write $w x$ for $u(w, x)$, and $w x^{n}$ for $u^{(n)}\left(w, x^{(n)}\right)$. For $w \in W$ and A Borel set in $X$, let $P_{1}(w, A)=P(w, A)$, and if $w \in W, m>1$ and $A$ Borel set in $X^{m}$, let the $m$-transfer probability of $w$ into $A$ be,

$$
P_{m}(w, A)=\int_{X} P\left(w, d x_{1}\right) \int_{X} P\left(w x_{1}, d x_{2}\right) \ldots \int_{X} P\left(w x^{(m-1)}, d x_{m}\right) d \chi_{A}\left(x^{(m)}\right.
$$

And for $w \in W, n, m \geq 1$ and $A \subset X^{m}$, define $P_{m}^{n}(w, A)=P_{n+m-1}\left(w, X^{n-1} \times A\right)$.
The following existence result was proved in [10].
Theorem 1.3. Let a random system with complete connections $\{(W, \mathcal{W}),(X, \mathcal{X}), u, P)$ and an arbitrary given point $w_{0} \in W$. Then, there exist a probability space $\left(\Omega, \mathcal{K}, P_{w_{0}}\right)$ and a sequence $\left(\xi_{n}\right)_{n \geq 1}$ of $X$-valued random variables defined on $\Omega$ such that, for all $m, n, q \geq 1$ and $A \in \mathcal{X}^{m}$, we have:
i) $P_{w_{0}}\left(\left[\xi_{n}, \ldots, \xi_{n+m-1}\right] \in A\right)=P_{m}^{n}\left(w_{0}, A\right)$,
ii) $P_{w_{0}}\left(\left[\xi_{n+q}, \ldots, \xi_{n+q+m-1}\right] \in A \mid \xi^{(n)}\right)=P_{m}^{q}\left(w_{0} \xi^{(n)}, A\right), P_{w_{0}}$-a.e.

Finally, the third main notion we will use is that of Smale skew product endomorphism, introduced in [17]. Let $I$ be a countable alphabet, and let $\Sigma_{I}^{+}$be the associated 1 -sided shift space, with shift map $\sigma: \Sigma_{I}^{+} \rightarrow \Sigma_{I}^{+}$. Given $\beta>0$, the metric $d_{\beta}$ on $\Sigma_{I}^{+}$is:

$$
d_{\beta}\left(\left(\omega_{n}\right)_{0}^{\infty},\left(\tau_{n}\right)_{0}^{\infty}\right)=\exp \left(-\beta \max \left\{n \geq 0:(0 \leq k \leq n) \Rightarrow \omega_{k}=\tau_{k}\right\}\right)
$$

with the standard convention that $e^{-\infty}=0$. All metrics $d_{\beta}, \beta>0$, on $\Sigma_{I}^{+}$are Hölder continuously equivalent and all induce the product topology on $\Sigma_{I}^{+}$. We define also the 2-sided shift space $\Sigma_{I}$ with the same metric as above.

For every $\omega \in \Sigma_{I}$ and integers $m \leq n$, define the ( $m, n$ )-truncation $\left.\omega\right|_{m} ^{n}=\omega_{m} \omega_{m+1} \ldots \omega_{n}$. Let $\Sigma_{I}^{*}$ be the set of finite words. For $\tau=\tau_{m} \tau_{m+1} \ldots \tau_{n}$, the cylinder from $m$ to $n$ is $[\tau]_{m}^{n}=\left\{\omega \in \Sigma_{I}:\left.\omega\right|_{m} ^{n}=\tau\right\}$. The family of cylinders from $m$ to $n$ is denoted by $C_{m}^{n}$. If $m=0$, write $[\tau]$ for $[\tau]_{m}^{n}$.

Let $\psi: \Sigma_{I} \rightarrow \mathbb{R}$ continuous. Topological pressure plays an important role in thermodynamic formalism and extends the notion of entropy (for eg [4], [22]). By extension, in our case the topological pressure of $\psi$ is: $\mathrm{P}(\psi)$ is,

$$
\begin{equation*}
\mathrm{P}(\psi):=\lim _{n \rightarrow \infty} \frac{1}{n} \log \sum_{\omega \in C_{0}^{n-1}} \exp \left(\sup \left(\left.S_{n} \psi\right|_{[\omega]}\right)\right) \tag{1.4}
\end{equation*}
$$

where the limit above exists by subadditivity. A shift-invariant Borel probability $\mu$ on the 2 -sided shift space $\Sigma_{I}$ with countable alphabet $I$, is called a Gibbs measure of $\psi$ if there exist constants $C \geq 1, P \in \mathbb{R}$ such that

$$
\begin{equation*}
C^{-1} \leq \frac{\mu\left(\left[\left.\omega\right|_{0} ^{n-1}\right]\right)}{\exp \left(S_{n} \psi(\omega)-P n\right)} \leq C \tag{1.5}
\end{equation*}
$$

for all $n \geq 1, \omega \in \Sigma_{I}$. From (1.5), if $\psi$ admits a Gibbs state, then $P=\mathrm{P}(\psi)$. The function $\psi: \Sigma_{I} \rightarrow \mathbb{R}$ is called summable if

$$
\sum_{e \in E} \exp \left(\sup \left(\left.\psi\right|_{[e]}\right)\right)<\infty
$$

In [17] we proved that a Hölder continuous $\psi: \Sigma_{I} \rightarrow \mathbb{R}$ is summable if and only if $\mathrm{P}(\psi)<\infty$, and that for every Hölder continuous summable function $\psi: \Sigma_{I} \rightarrow \mathbb{R}$ there exists a unique Gibbs state $\mu_{\psi}$ on $\Sigma_{I}$, and the measure $\mu_{\psi}$ is ergodic. We also showed that if $\psi: \Sigma_{I} \rightarrow \mathbb{R}$ is a Hölder continuous summable function, then

$$
\sup \left\{\mathrm{h}_{\mu}(\sigma)+\int_{\Sigma_{\mathrm{I}}} \psi \mathrm{~d} \mu: \mu \circ \sigma^{-1}=\mu \text { and } \int \psi \mathrm{d} \mu>-\infty\right\}=\mathrm{P}(\psi)=\mathrm{h}_{\mu_{\psi}}(\sigma)+\int_{\Sigma_{I}} \psi d \mu_{\psi}
$$

and this supremum is attained only at $\mu_{\psi}$. Let now the measurable partition of $\Sigma_{I}$,

$$
\begin{equation*}
\mathcal{P}_{-}=\left\{\left[\left.\omega\right|_{0} ^{+\infty}\right]: \omega \in \Sigma_{I}\right\}=\left\{[\omega]: \omega \in \Sigma_{I}^{+}\right\} \tag{1.6}
\end{equation*}
$$

If $\mu$ is a probability on $\Sigma_{I}$, let the canonical conditional measures associated to $\mathcal{P}_{-}([21])$,

$$
\left\{\bar{\mu}^{\tau}: \tau \in \Sigma_{I}^{+}\right\}
$$

Then $\bar{\mu}^{\tau}$ is a probability measure on the cylinder $\left[\left.\tau\right|_{0} ^{+\infty}\right]$ and we also write $\bar{\mu}^{\omega}, \omega \in \Sigma_{I}^{+}$, for the conditional measure on $[\omega]$. The canonical projection (truncation) is:

$$
\pi_{0}: \Sigma_{I} \rightarrow \Sigma_{I}^{+}, \pi_{0}(\tau)=\left.\tau\right|_{0} ^{\infty}, \tau \in \Sigma_{I}
$$

Recall that the system $\left\{\bar{\mu}^{\omega}: \omega \in \Sigma_{I}^{+}\right\}$of conditional measures is uniquely determined (up to measure zero), by: $\int_{\Sigma_{I}} g d \mu=\int_{\Sigma_{I}^{+}} \int_{[\omega]} g d \bar{\mu}^{\omega} d\left(\mu \circ \pi_{0}^{-1}\right)(\omega), \forall g \in L^{1}(\mu)([21])$.

We introduced the notion of Smale skew product endomorphisms.
Definition 1.4. [17] Let $(Y, d)$ be a complete bounded metric space. For every $\omega \in \Sigma_{I}^{+}$let $Y_{\omega} \subset Y$ be an arbitrary set and let $T_{\omega}: Y_{\omega} \longrightarrow Y_{\sigma(\omega)}$ be a continuous injective map. Let $\hat{Y}:=\bigcup_{\omega \in \Sigma_{I}^{+}}\{\omega\} \times Y_{\omega} \subset \Sigma_{I}^{+} \times Y$, and define the map $T: \hat{Y} \longrightarrow \hat{Y}, T(\omega, y)=\left(\sigma(\omega), T_{\omega}(y)\right)$. The pair $(\hat{Y}, T: \hat{Y} \rightarrow \hat{Y})$ is called a model Smale endomorphism if there exists $\lambda>1$ such that for all $\omega \in \Sigma_{I}^{+}$and all $y_{1}, y_{2} \in Y_{\omega}, d\left(T_{\omega}\left(y_{2}\right), T_{\omega}\left(y_{1}\right)\right) \leq \lambda^{-1} d\left(y_{2}, y_{1}\right)$.

$$
\begin{gathered}
\text { If } \tau=\left(\tau_{-n}, \ldots, \tau_{0}, \tau_{1}, \ldots\right) \text { let } T_{\tau}^{n}=T_{\tau| |_{-1}^{\infty}} \circ T_{\tau \mid{ }_{-2}^{\infty}} \circ \ldots \circ T_{\tau| |_{n}^{\infty}}: Y_{\tau} \longrightarrow Y_{\left.\tau\right|_{0} ^{\infty}} . \text { If } \tau \in \Sigma_{I} \text { let, } \\
T_{\tau}^{n}:=T_{\tau \mid{ }_{-n}}^{n}:=T_{\tau \mid{ }_{-1}^{\infty}} \circ T_{\tau| |_{-2}^{\infty}} \circ \ldots \circ T_{\left.\tau\right|_{-n} ^{\infty}}: Y_{\tau \mid{ }_{-n}^{\infty}} \longrightarrow Y_{\left.\tau\right|_{0} ^{\infty}}
\end{gathered}
$$

Then the sets $\left(T_{\tau}^{n}\left(Y_{\tau \mid \propto_{n}}\right)\right)_{n=0}^{\infty}$ form a descending sequence, and $\operatorname{diam}\left(\overline{T_{\tau}^{n}\left(Y_{\tau \mid{ }_{-}}\right)}\right) \leq \lambda^{-n} \operatorname{diam}(Y)$. As $(Y, d)$ is complete, $\bigcap_{n=1}^{\infty} \overline{T_{\tau}^{n}\left(Y_{\tau \mid-n}\right)}$ is a point denoted by $\hat{\pi}_{2}(\tau)$, which defines the map

$$
\begin{equation*}
\hat{\pi}_{2}: \Sigma_{I} \longrightarrow Y \tag{1.7}
\end{equation*}
$$

and define also $\hat{\pi}: \Sigma_{I} \rightarrow \Sigma_{I}^{+} \times Y$ by

$$
\begin{equation*}
\hat{\pi}(\tau)=\left(\left.\tau\right|_{0} ^{\infty}, \hat{\pi}_{2}(\tau)\right) \tag{1.8}
\end{equation*}
$$

and the truncation to non-negative indices by $\pi_{0}: \Sigma_{I} \longrightarrow \Sigma_{I}^{+}, \quad \pi_{0}(\tau)=\left.\tau\right|_{0} ^{\infty}$.
Now assume $Y_{\omega}=Y, \forall \omega \in \Sigma_{I}^{+}$, and for an arbitrary $\omega \in \Sigma_{I}^{+}$denote the $\hat{\pi}_{2}$-projection of the cylinder $[\omega] \subset \Sigma_{I}, J_{\omega}:=\hat{\pi}_{2}([\omega]) \subset Y$, and call these sets the stable Smale fibers of $T$. The global invariant set,

$$
J:=\hat{\pi}\left(\Sigma_{I}\right)=\bigcup_{\omega \in \Sigma_{I}^{+}}\{\omega\} \times J_{\omega} \subset \Sigma_{I}^{+} \times Y
$$

is called the Smale space induced by $T$. Then the skew-product system

$$
\begin{equation*}
T: J \longrightarrow J \tag{1.9}
\end{equation*}
$$

is called the Smale endomorphism generated by $T: \hat{Y} \longrightarrow \hat{Y}$.
Now suppose more conditions about $Y_{\omega}, \omega \in \Sigma_{I}^{+}$and the maps $T_{\omega}: Y_{\omega} \rightarrow Y_{\sigma(\omega)}$, namely:
(a) $Y_{\omega}$ is a closed bounded subset of $\mathbb{R}^{d}$, with some $d \geq 1$ such that $\overline{\operatorname{Int}\left(Y_{\omega}\right)}=Y_{\omega}$.
(b) Each map $T_{\omega}: Y_{\omega} \rightarrow Y_{\sigma(\omega)}$ extends to a $C^{1}$ conformal embedding from $Y_{\omega}^{*}$ to $Y_{\sigma(\omega)}^{*}$, where $Y_{\omega}^{*}$ is a bounded connected open subset of $\mathbb{R}^{d}$ containing $Y_{\omega}$. Then $T_{\omega}$ denotes also this extension and assume that the maps $T_{\omega}: Y_{\omega}^{*} \rightarrow Y_{\sigma(\omega)}^{*}$ satisfy:
(c) There is $\lambda>1$ so that $d\left(T_{\omega}\left(y_{1}\right), T_{\omega}\left(y_{2}\right)\right) \leq \lambda^{-1} d\left(y_{1}, y_{2}\right), \forall \omega \in \Sigma_{I}^{+}, y_{1}, y_{2} \in Y_{\omega}^{*}$.
(d) (Bounded Distortion Property 1) There are constants $\alpha>0, H>0$ s.t $\forall y, z \in Y_{\omega}^{*}$,

$$
|\log | T_{\omega}^{\prime}(y)|-\log | T_{\omega}^{\prime}(z)| | \leq H\|y-z\|^{\alpha}
$$

(e) The function $\Sigma_{I} \ni \tau \longmapsto \log \left|T_{\omega}\left(\hat{\pi}_{2}(\eta)\right)\right| \in \mathbb{R}$ is Hölder continuous, where $\omega=\pi_{0}(\tau)$.
(f) (Open Set Condition) For every $\omega \in \Sigma_{I}^{+}$and for all $a, b \in I$ with $a \neq b$, we have $T_{a \omega}\left(\operatorname{Int}\left(Y_{a \omega}\right)\right) \cap T_{b \omega}\left(\operatorname{Int}\left(Y_{b \omega}\right)\right)=\emptyset$.
(g) (Strong Open Set Condition) There exists a measurable function $\delta: \Sigma_{I}^{+} \rightarrow(0, \infty)$ such that for every $\omega \in \Sigma_{I}^{+}, J_{\omega} \cap\left(Y_{\omega} \backslash \bar{B}\left(Y_{\omega}^{c}, \delta(\omega)\right) \neq \emptyset\right.$.
A skew product Smale endomorphism satisfying conditions (a)-(g) will be called in the sequel a conformal Smale endomorphism.

We see that a countable IFS with place-dependent probabilities is a particular case of random system with complete connections with state space $V$, index space $=$ countable alphabet $I$, and probability transition function determined by $p_{i}(\cdot), i \in I$, namely

$$
\begin{equation*}
P(x, B)=\sum_{i \in I} p_{i}(x) \delta_{\phi_{i}(x)}(B) \tag{1.10}
\end{equation*}
$$

Also, to a conformal Smale endomorphism $T$ with disjoint fibers $J_{\omega}, \omega \in \Sigma_{I}^{+}$, we will associate a random system with complete connections in the next result. Recall that $\pi_{0}: \Sigma_{I} \rightarrow \Sigma_{I}^{+}, \pi_{0}(\eta)=\left.\eta\right|_{0} ^{\infty}, \eta \in \Sigma_{I}$.

Theorem 1.5. Let the conformal Smale endomorphism $T: J \rightarrow J$ defined in (1.9), where we assume the spaces $Y_{\omega}=Y \subset \mathbb{R}^{d}, \forall \omega \in \Sigma_{I}^{+}$. Consider also a shift invariant measure $\mu$ on $\Sigma_{I}$. Define the quadruple $((W, \mathcal{W}),(X, \mathcal{X}), u, P)$ by:
a) State space $W=\left\{(\omega, x), x \in J_{\omega}, \omega \in \Sigma_{I}^{+}\right\}$with the Borel $\sigma$-algebra $\mathcal{W}$ induced from the Borel $\sigma$-algebra of the product space $\Sigma_{I}^{+} \times Y \subset \Sigma_{I}^{+} \times \mathbb{R}^{d}$, and index space $X=\Sigma_{I}^{+}$with its Borel $\sigma$-algebra $\mathcal{X}$;
b) For $\tau \in \Sigma_{I}^{+}$and $w \in W, w=(\omega, x), x \in J_{\omega}$, define the map $u(w, \tau)=\left(\sigma \omega, T_{\omega}(x)\right)$;
c) For those $\omega \in \Sigma_{I}^{+}$for which the conditional measure $\bar{\mu}^{\omega}$ is defined (their set has $\pi_{0 *} \mu$ measure equal to 1 ), and for $w=(\omega, x), x \in J_{\omega}$, define the probability transition function $P_{w}(\cdot)$ by $P_{w}:=\bar{\mu}^{\omega}$.

Then, $((W, \mathcal{W}),(X, \mathcal{X}), u, P)$ is a random system with complete connections.
Proof. On $W$ we take the $\sigma$-algebra of Borel sets induced from $\Sigma_{I}^{+} \times \mathbb{R}^{d}$. We defined the map $u(w, \tau)=\left(\sigma \omega, T_{\omega}(x)\right)$, for $w=(\omega, x) \in W, \omega \in \Sigma_{I}^{+}$and $x \in J_{\omega}$. On the other hand recall that $J_{\omega}$ is the set of points of the form $\hat{\pi}_{2}(\eta)$ for $\eta \in[\omega] \subset \Sigma_{I}$. So from (1.7) the map $u(\cdot, \cdot)$ is well-defined, since if $x=\hat{\pi}_{2}(\eta)=T_{\eta_{-1} \omega} \circ T_{\eta_{-2} \eta_{-1} \omega} \circ \ldots \in J_{\omega}$, then

$$
T_{\omega}(x)=\hat{\pi}_{2}(\sigma \eta) \in J_{\sigma \omega}
$$

as $\sigma \eta \in[\sigma \omega]$. Then, we use condition (e) from the definition of conformal Smale endomorphisms, to obtain that $u: W \times X \rightarrow W$ is measurable.

Next, for $\pi_{0 *} \mu$-a.e $\omega \in \Sigma_{I}^{+}$, the conditional measure $\bar{\mu}^{\omega}$ is defined on the cylinder $[\omega] \subset \Sigma_{I}$, and this cylinder can be identified with $\Sigma_{I}^{+}$. If $A$ is a Borel set in $\Sigma_{I}^{+}$and $w=(\omega, x) \in$ $W, x \in J_{\omega}$, define $P(w, A)=\bar{\mu}^{\omega}(A)$, thus $P_{w}$ can be viewed as a probability measure on $\Sigma_{I}^{+}$. From the uniqueness of the system of conditional measures associated to the partition $\mathcal{P}_{-}$(see [21]) with the property that

$$
\int_{\Sigma_{I}} g(\xi) d \mu=\int_{\Sigma_{I}^{+}} \int_{[\omega]} g(\xi) d \bar{\mu}^{\omega}(\xi) d \pi_{0 *} \mu(\omega),
$$

for any integrable function $g: \Sigma_{I} \rightarrow \mathbb{R}$, we obtain that for any set $A$ as above, $P_{w}(A)$ depends measurably on $\omega \in \Sigma_{I}^{+}$, where $w=(\omega, x), x \in J_{\omega}$. Hence the function

$$
P(\cdot, A): W \rightarrow \mathbb{R}, w \rightarrow P(w, A)
$$

is also measurable with respect to the $\sigma$-algebra $\mathcal{W}$ induced on $W$ from $\Sigma_{I}^{+} \times \mathbb{R}^{d}$.

## 2. Stationary measures of countable systems with overlaps and PLACE-DEPENDENT PROBABILITIES

In this section we study the case of a countable IFS with place-dependent probabilities and overlaps, and prove the exact dimensionality of stationary measures, and compute the pointwise (and Hausdorff) dimension.

Consider a system of smooth contractions defined on a compact set $V \subset \mathbb{R}^{D}$ indexed by a countable alphabet $I, \mathcal{S}=\left\{\phi_{i}: V \longrightarrow V\right\}_{i \in I}$ with limit set $\Lambda$, and the weights $p_{i}: V \rightarrow \mathbb{R}$, $i \in I$, satisfying for any $x \in V$,

$$
\begin{equation*}
\sum_{i \in I} p_{i}(x)=1 \tag{2.1}
\end{equation*}
$$

$\Sigma_{I}^{+}$denotes the 1-sided shift space with alphabet $I$. If $i_{1}, \ldots, i_{n} \in I, n \geq 1$, denote

$$
\phi_{i_{1} \ldots i_{n}}:=\phi_{i_{1}} \circ \ldots \circ \phi_{i_{n}}
$$

Let $\pi: \Sigma_{I}^{+} \rightarrow \Lambda, \pi(\omega)=\lim _{n \rightarrow \infty} \phi_{i_{1} i_{2} \ldots i_{n}}$ if $\omega=\left(i_{1}, i_{2}, \ldots\right) \in \Sigma_{I}^{+}$, be the canonical coding map for the limit set $\Lambda$.

Assume also that $p_{i}(\cdot)$ depend uniformly Hölder continuously on $x \in V$, for $i \in I$, i.e. there exist constants $\alpha, C>0$ such that for all $i \in I$ and all $x, y \in V$,

$$
\begin{equation*}
\left|p_{i}(x)-p_{i}(y)\right| \leq C|x-y|^{\alpha} \tag{2.2}
\end{equation*}
$$

The transfer probability in this case is $P(x, B):=\sum_{i \in I} p_{i}(x) \delta_{\phi_{i}(x)}(B)$, and the transfer operator $\mathcal{L}: \mathcal{C}(V) \rightarrow \mathcal{C}(V)$ is given by:

$$
\mathcal{L}(f)(x)=\int_{X} f(y) P(x, d y)
$$

A measure $\mu$ on $V$ is called stationary if it is a fixed point of the dual operator of $\mathcal{L}$,

$$
\mathcal{L}^{*}(\nu)(B)=\int P(x, B) d \nu(x)=\sum_{i \in I} \int_{\phi_{i}^{-1}(B)} p_{i}(x) d \nu(x) .
$$

Define also the Lyapunov exponent of a shift-invariant measure $\mu$ on $\Sigma_{I}^{+}$by:

$$
\chi_{\mu}:=-\int_{\Sigma_{I}^{+}} \log \left|\phi_{\omega_{1}}^{\prime}(\pi(\sigma \omega))\right| d \mu(\omega) .
$$

Let us now recall the notion of projection entropy for an invariant measure of a countable iterated function system from [16]; this is a generalization of the notion of projection entropy from the finite case of [7]. As a matter of fact, in [16] we defined the projection entropy in the more general case of random countable iterated function systems, but here we need it only for countable deterministic systems.

So let $\mathcal{S}=\left\{\phi_{i}, i \in I\right\}$ be a countable system and $\mu$ be a $\sigma$-invariant probability measure on $\Sigma_{I}^{+}$. Denote by $\xi$ the partition of $\Sigma_{I}^{+}$into initial 1 -cylinders, and by $\epsilon_{\mathbb{R}^{D}}$ the point partition of $\mathbb{R}^{D}$, and by $\pi: \Sigma_{I}^{+} \rightarrow \Lambda$ the canonical coding map for the limit set $\Lambda \subset \mathbb{R}^{D}$ of the function system $\mathcal{S}$. Then $\pi^{-1} \epsilon_{\mathbb{R}^{D}}$ and $\sigma^{-1}\left(\pi^{-1} \epsilon_{\mathbb{R}^{D}}\right)$ are measurable partitions of $\Sigma_{I}^{+}$. The projection entropy of $\mu$ with respect to $\mathcal{S}$ is defined then by,

$$
\begin{equation*}
h_{\mu}(\mathcal{S}):=H_{\mu}\left(\xi \mid \sigma^{-1}\left(\pi^{-1} \epsilon_{\mathbb{R}^{D}}\right)\right)-H_{\mu}\left(\xi \mid \pi^{-1} \epsilon_{\mathbb{R}^{D}}\right) . \tag{2.3}
\end{equation*}
$$

We prove next that the stationary measure from Theorem 1.3 is exact dimensional.

Theorem 2.1. In the above setting if the system $\mathcal{S}=\left\{\phi_{i}, i \in I\right\}$ is countable and conformal and if the probabilities $\left\{p_{i}(\cdot), i \in I\right\}$ satisfy (2.1)-(2.2), then the stationary measure $\tilde{\mu}_{P}$ for the system $\mathcal{S}$ with place-dependent probabilities $P=\left\{p_{i}(\cdot), i \in I\right\}$ is exact dimensional, and

$$
H D\left(\tilde{\mu}_{P}\right) \leq \frac{h_{\mu_{\psi}}(\sigma)}{\chi_{\mu_{\psi}}}
$$

where $\psi: \Sigma_{I}^{+} \rightarrow \mathbb{R}, \psi(\omega):=\log p_{\omega_{0}}(\pi(\sigma \omega)), \omega \in \Sigma_{I}^{+}$, and $\mu_{\psi}$ is the equilibrium measure of $\psi$ on $\Sigma_{I}^{+}$.

Proof. First let us define the potential $\psi: \Sigma_{I}^{+} \rightarrow \mathbb{R}$ by:

$$
\psi(\omega):=\log p_{\omega_{0}}(\pi(\sigma \omega)), \omega \in \Sigma_{I}^{+}
$$

From (2.1) and (2.2) it follows that $\psi$ is summable and Hölder continuous on $\Sigma_{I}^{+}$. Then there exists an equilibrium measure $\mu_{\psi}$ on $\Sigma_{I}^{+}$, which projects to the probability measure $\nu_{p}$ on $\Lambda$, associated to the system of weights $p:=\left(p_{i}, i \in I\right)$. The transfer operator $\mathcal{L}: \mathcal{C}\left(\Sigma_{I}^{+}\right) \rightarrow C\left(\Sigma_{I}^{+}\right)$is,

$$
\mathcal{L}(\phi)(\omega):=\sum_{i=1}^{\infty} p_{i}(\pi \omega) \phi(i \omega)
$$

where $\mathcal{C}\left(\Sigma_{I}^{+}\right)$is the space of bounded continuous real-valued functions on $\Sigma_{I}^{+}$. We see from above that

$$
\mathcal{L}(\phi)(\omega)=\sum_{i=1}^{\infty} e^{\psi(i \omega)} \phi(i \omega)
$$

For such transfer operators, it was proved (see [11]) that if $\psi$ is Hölder continuous and summable, then there exists a unique probability measure $\tilde{\nu}_{\psi}$ on $\Sigma_{I}^{+}$, such that

$$
\mathcal{L}^{*}\left(\tilde{\nu}_{\psi}\right)=\tilde{\nu}_{\psi}
$$

The projection, denoted by $\tilde{\mu}_{P}$, of the measure $\tilde{\nu}_{\psi}$ onto the limit set $\Lambda$ of $\mathcal{S}$, is the stationary measure of the system $\mathcal{S}$ with the place-dependent probabilities $P$. Hence,

$$
\tilde{\mu}_{P}=\pi_{*} \tilde{\nu}_{\psi}
$$

On the other hand, for the expanding map $\sigma: \Sigma_{I}^{+} \rightarrow \Sigma_{I}^{+}$and the Hölder continuous potential $\psi$, let us notice that there exists also a shift-invariant equilibrium measure $\mu_{\psi}$ of $\psi$ on $\Sigma_{I}^{+}$, and moreover there exists a function $\theta$ such that

$$
\theta \tilde{\mu}_{\psi}=\mu_{\psi}
$$

Moreover this function $\theta$ satisfies

$$
\theta \geq M>0
$$

for some constant $M$ depending on $\psi$ (see [11]). On the other hand, the projection of the shift-invariant equilibrium measure $\mu_{\psi}$ onto the limit set $\Lambda$ is denoted by $\mu_{P}$ and we have

$$
\mu_{P}=\pi_{*} \mu_{\psi}
$$

In [16] we proved that the projection $\mu_{P}$ of the invariant measure $\mu_{\psi}$ is exact dimensional even when the system $\mathcal{S}$ has overlaps. And we found a formula for its pointwise dimension, involving the projection entropy (recalled above). Indeed in the main Theorem of [16], it is enough to consider a random system where the parameter space consists of only one point,
and to take the identity as the evolution map on the space of parameters. This means that for $\mu_{P}$-a.e $x \in \Lambda$,

$$
\begin{equation*}
\lim _{r \rightarrow 0} \frac{\log \mu_{P}(B(x, r))}{\log r}=\delta=\frac{h_{\mu_{\psi}}(\mathcal{S})}{\chi_{\mu_{\psi}}} \tag{2.4}
\end{equation*}
$$

and $\delta$ does not depend on $x \in \Lambda$. But $\mu_{P}(B(x, r))=\mu_{\psi}\left(\pi^{-1}(B(x, r))\right)$, and we know that $\mu_{\psi}=\theta \tilde{\mu}_{\psi}$, hence

$$
\tilde{\mu}_{\psi}\left(\pi^{-1}(B(x, r))\right)=\int_{\pi^{-1}(B(x, r))} \theta d \mu_{\psi}
$$

On the other hand let us recall that $\theta>M$ and that $\theta$ is a continuous bounded function on $\Sigma_{I}^{+}$. Hence using (2.4) and the fact that $\tilde{\mu}_{P}=\pi_{*} \tilde{\mu}_{\psi}$, we see that for $\mu_{P}$-a.e $x \in \Lambda$,

$$
\begin{equation*}
\lim _{r \rightarrow 0} \frac{\log \tilde{\mu}_{P}(B(x, r))}{\log r}=\delta \tag{2.5}
\end{equation*}
$$

Therefore, the stationary measure $\tilde{\mu}_{p}$ of the system $\mathcal{S}$ with the place-dependent probabilities $P$ is exact dimensional, and from (2.5) its Hausdorff (and pointwise) dimension is given by:

$$
\begin{equation*}
H D\left(\tilde{\mu}_{p}\right)=\frac{h_{\mu_{\psi}}(\mathcal{S})}{\chi_{\mu_{\psi}}} \tag{2.6}
\end{equation*}
$$

where $h_{\mu_{\psi}}(\mathcal{S})$ is the projection entropy of $\mu_{\psi}$ with respect to $\mathcal{S}$ and $\chi_{\mu_{\psi}}$ is its Lyapunov exponent. From the definition (2.3) of $h_{\mu_{\psi}}(\mathcal{S})$ it follows that $h_{\mu_{\psi}}(\mathcal{S}) \leq h_{\mu_{\psi}}(\sigma)$. Therefore, we obtain the desired dimension formula.

## 3. Families of fractals and unfoldings of countable IFS.

We want now to associate a Smale endomorphism to a countable iterated function systems with overlaps, by unfolding, in such a way as to control the structure of overlappings. We will consider in general equilibrium measures of real-valued summable functions $\psi$ on $\Sigma_{I}$.

Let us consider a countable conformal IFS with overlaps $\mathcal{S}=\left\{\phi_{i}, i \in I\right\}$, where the maps $\phi_{i}: V \rightarrow V$ are contractions defined on a neighbourhood of a compact set $V \subset \mathbb{R}^{D}$, and $\left|\phi_{i}^{\prime}\right|<\alpha<1, i \in I$ on $V$. Denote by $\Lambda$ the limit set of $\mathcal{S}$, and assume that $\Lambda$ is not contained in the boundary of $V$. Since we work with a countable system, the limit set $\Lambda$ may be non-compact. Assume that $I=\mathbb{N}^{*}$ and that $\mathcal{S}$ satisfies the Bounded Distortion Property, i.e there exist constants $H>0, \beta>0$ such that for all $i \in I$,

$$
\begin{equation*}
|\log | \phi_{i}^{\prime}(y)|-\log | \phi_{i}^{\prime}(z)| | \leq H|y-z|^{\beta}, \forall y, z \in V \tag{3.1}
\end{equation*}
$$

We will associate to $\mathcal{S}$, a Smale skew product $T$ and a random system with complete connections, with the goal to separate the images of the compositions of maps $\phi_{i}$ along any given sequence $\omega=\left(\omega_{1}, \omega_{2}, \ldots\right) \in \Sigma_{I}^{+}$. This is realised by an inductive process of unfolding the overlaps of $\mathcal{S}$. Recall that $\Lambda$ is the limit set of $\mathcal{S}$.

Assume that for any point $x \in \Lambda$, the $\mathcal{S}$-images $\phi_{i}(x), i \in I$ of $x$, satisfy the following Non-accumulation Condition:

$$
\begin{equation*}
\phi_{i}(x) \notin \overline{\left\{\phi_{j}(x), j \in I \backslash\{i\}\right\}}, \forall i \in I \tag{3.2}
\end{equation*}
$$

This condition is quite general, and it can be checked on many systems (see for instance the examples in [18]). Recall that $\pi(\omega)=\phi_{\omega_{1} \omega_{2} \ldots}$ is the canonical projection from $\Sigma_{I}^{+}$to the limit set $\Lambda$. Then, for an arbitrary $\omega=\left(\omega_{1}, \omega_{2}, \ldots\right) \in \Sigma_{I}^{+}$, we define inductively the contractions $T_{i \omega}$ for $i \in I$. Let us start by defining

$$
T_{1 \omega}:=\phi_{1 \omega_{1} \ldots \omega_{n_{1}(\omega)}}=\phi_{1} \circ \phi_{\omega_{1}} \circ \ldots \circ \phi_{\omega_{n_{1}(\omega)}},
$$

where $n_{1}(\omega)$ is defined as the smallest integer $n_{1} \geq 1$ such that

$$
\phi_{j}(\pi(\omega)) \notin \phi_{1 \omega_{1} \ldots \omega_{n_{1}}}(V), \text { for all } j \neq 1
$$

Next, since from above $\phi_{2}(\pi(\omega)) \notin \phi_{1 \omega_{1} \ldots \omega_{n_{1}}}(V)$, take $n_{2}(\omega)$ to be the smallest integer $n_{2}>n_{1}$ such that

$$
\phi_{2 \omega_{1} \ldots \omega_{n_{2}(\omega)}}(V) \cap \phi_{1 \omega_{1} \ldots \omega_{n_{1}(\omega)}}(V)=\emptyset, \text { and } \phi_{j}(\pi \omega) \notin \phi_{2 \omega_{1} \ldots \omega_{n_{2}(\omega)}}(V), \text { for } j \neq 2
$$

Then we define

$$
T_{2 \omega}:=\phi_{2 \omega_{1} \ldots \omega_{n_{2}(\omega)}}
$$

Inductively, if we defined $n_{k}(\omega)$ up to some $k \geq 1$, we now define $n_{k+1}(\omega) \geq 1$ as the smallest integer $n_{k+1}>n_{k}$ with the property that
$\phi_{(k+1) \omega_{1} \ldots \omega_{n_{k+1}(\omega)}}(V) \cap \phi_{j \omega_{1} \ldots \omega_{n_{j}(\omega)}}(V)=\emptyset, 1 \leq j \leq k$, and $\phi_{\ell}(\pi \omega) \notin \phi_{(k+1) \omega_{1} \ldots \omega_{n_{k+1}(\omega)}}(V), \ell \neq k+1$, and then the fiber map

$$
\begin{equation*}
T_{k+1 \omega}:=\phi_{(k+1) \omega_{1} \ldots \omega_{n_{k+1}(\omega)}} \tag{3.4}
\end{equation*}
$$

Thus we constructed for any $\eta \in \Sigma_{I}^{+}$, a contraction $T_{\eta}=T_{\eta_{1} \sigma(\eta)}$ as above. Consider now a sequence $\tau \in \Sigma_{I}$, then for any $n \geq 1$ define the map

$$
T_{\tau}^{n}=T_{\tau_{-1}^{\infty}} \circ T_{\tau \mid{ }_{-2}} \circ \ldots \circ T_{\tau \mid{ }_{-n}}
$$

Then $\hat{\pi}_{2}(\tau)=\bigcap_{n=1}^{\infty} \overline{T_{\tau}^{n}(V)}$, and the fractal $J_{\omega}=\hat{\pi}_{2}([\omega])$ is contained in $\Lambda$, where $[\omega]$ is the cylinder in $\Sigma_{I}$ determined by $\omega \in \Sigma_{I}^{+}$.

Definition 3.1. Let us define the space $\hat{Y}:=\Sigma_{I}^{+} \times V$, and the skew product $T: \hat{Y} \rightarrow \hat{Y}$,

$$
T(\omega, x)=\left(\sigma(\omega), T_{\omega}(x)\right),(\omega, x) \in \hat{Y}
$$

We will call $T: \Sigma_{I}^{+} \times V \rightarrow \Sigma_{I}^{+} \times V$ the maximal Smale system associated to the countable IFS with overlaps $\mathcal{S}$.

From definition we see that the map $T_{\omega}$ depends on the whole sequence $\omega \in \Sigma_{I}^{+}$, and not just on the projection point $\pi(\omega)$. Thus the dynamical system $(\hat{Y}, T)$ describes how the overlappings are formed through iterations.

If $\mu$ is a probability measure on $\Sigma_{I}$, define its Lyapunov exponent with respect to $T$ as

$$
\chi_{\mu}(\sigma)=-\int_{\Sigma_{I}} \log \left|T_{\left.\tau\right|_{0} ^{\infty}}^{\prime}\left(\hat{\pi}_{2}(\tau)\right)\right| d \mu(\tau)
$$

Denote by $\pi_{0}: \Sigma_{I} \rightarrow \Sigma_{I}^{+}, \pi_{0}(\tau)=\left(\tau_{0}, \tau_{1}, \ldots\right)$, the canonical truncation map.

We constructed above the family of random fractals $J_{\omega} \subset \Lambda, \omega \in \Sigma_{I}^{+}$, and now we prove that certain invariant probability measures supported on $J_{\omega}$ are exact dimensional.

Theorem 3.2. Let a countable IFS with overlaps $\mathcal{S}$ as above satisfying (3.2), and let $T$ be its associated maximal Smale system. Consider $\psi: \Sigma_{I} \rightarrow \mathbb{R}$ summable Hölder continuous function with equilibrium measure $\mu_{\psi}$, and let $\nu_{\psi}$ be the canonical projection $\pi_{0, *} \mu_{\psi}$ of $\mu_{\psi}$ on $\Sigma_{I}^{+}$. Take the conditional measure $\mu_{\psi}^{\omega}$ of $\mu_{\psi}$ on the cylinder $[\omega]$ and let $\nu_{\psi}^{\omega}:=\hat{\pi}_{2 *} \mu_{\psi}^{\omega}$ be its projection on $J_{\omega}$, with $\hat{\pi}_{2}$ defined in 1.8. Then, for $\nu_{\psi}$-a.e $\omega \in \Sigma_{I}^{+}$, the measure $\nu_{\psi}^{\omega}$ is exact dimensional on the sub-fractal $J_{\omega} \subset \Lambda$, and

$$
H D\left(\nu_{\psi}^{\omega}\right)=\frac{h_{\mu_{\psi}}(\sigma)}{\chi_{\mu_{\psi}}(\sigma)}
$$

Proof. First let us notice that, from our construction, for any $i \neq j$,

$$
T_{i \omega}(V) \cap T_{j \omega}(V)=\emptyset
$$

Therefore the open set condition in fibers from the definition of the Smale skew-product, is satisfied. Moreover if $\left\|\phi_{i}^{\prime}\right\|<\alpha<1, i \in I$, it follows that the same uniform contractivity condition is satisfied by all the maps $T_{\omega}, \omega \in \Sigma_{I}^{+}$. Hence the uniform contractivity of the maps $T_{\omega}$ is satisfied.

Let us see now if the maps $T_{\omega}$ satisfy Bounded Distortion Property (BDP). For this consider an arbitrary $\omega \in \Sigma_{I}^{+}$. Then $T_{\omega}=\phi_{\omega_{1} \omega_{2} \ldots \omega_{n}}$ for some integer $n$ which depends on $\omega$. Thus, there exists $L \in(0,1)$ and $H^{\prime}>0$ such that for any $y, z \in V$, we have

$$
\begin{aligned}
& |\log | T_{\omega}^{\prime}(y)|-\log | T_{\omega}^{\prime}(z)| | \leq \\
& \leq|\log | \phi_{\omega_{1}}\left(\phi_{\omega_{2} \ldots \omega_{n}}(y)\right)|-\log | \phi_{\omega_{1}}\left(\phi_{\omega_{2} \ldots \omega_{n}}(z)\right)| |+\ldots+|\log | \phi_{\omega_{n}}^{\prime}(y)|-\log | \phi_{\omega_{n}}^{\prime}(z)| | \leq \\
& \leq H|y-z|^{\beta}\left(1+L+\ldots+L^{n}\right)=H^{\prime}|y-z|^{\beta}
\end{aligned}
$$

Another condition in the definition of a conformal Smale skew product is the Holder continuity of the real-valued map on $\Sigma_{I}$ given by:

$$
\begin{equation*}
\tau \longrightarrow \log \left|T_{\tau}^{\prime}\left(\hat{\pi}_{2}(\tau)\right)\right| \tag{3.5}
\end{equation*}
$$

Let us take $\omega \in \Sigma_{I}^{+}$and $\tau \in \Sigma_{I}$ such that $\tau \in[\omega]$. Then $T_{\tau}=T_{\tau_{-1} \omega}$, and consider the integer $n_{\tau_{-1}}(\omega)$. Then from definition, we have

$$
T_{\tau_{-1} \omega}=\phi_{\tau_{-1}} \circ \phi_{\omega_{1}} \circ \ldots \circ \phi_{\omega_{n_{\tau_{-1}}(\omega)}}
$$

But if $\eta \in\left[\tau_{-m} \ldots \tau_{m}\right]$ for $m \geq n_{\tau_{-1}(\omega)}$, we have that $T_{\tau}=T_{\eta}$. On the other hand, due to the form of $\hat{\pi}_{2}(\tau)$ given in (1.7), we have

$$
d\left(\hat{\pi}_{2}(\tau), \hat{\pi}_{2}(\eta)\right) \leq C \frac{1}{2^{m}}
$$

for some constant $C$ independent of $\tau, \eta, m$. Hence the map from (3.5) is indeed Hölder continuous on $\Sigma_{I}$.

Therefore, the maximal Smale system $T: \hat{Y} \rightarrow \hat{Y}$ defined above, where we defined

$$
T_{\left.\tau\right|_{0} ^{\infty}}=\phi_{\tau_{0} \tau_{1} \ldots \tau_{n\left(\left.\sigma \tau\right|_{0} ^{\infty}\right)}},
$$

satisfies the properties of a conformal Smale skew product endomorphism. We then apply our result from [17], to obtain the exact dimensionality and the formula for the dimension of the projection measures $\nu_{\psi}^{\omega}$ on $J_{\omega} \subset \Lambda$, for $\nu_{\psi^{-}}$-a.e. $\omega \in \Sigma_{I}^{+}$. In particular, it follows that

$$
H D\left(\nu_{\psi}^{\omega}\right)=\frac{h_{\mu_{\psi}}(\sigma)}{\chi_{\mu_{\psi}}(\sigma)}
$$

where $h_{\mu_{\psi}}(\sigma)$ is the entropy of the measure $\mu_{\psi}$, and $\chi_{\mu_{\psi}}(\sigma)$ is its Lyapunov exponent.

Examples of functions $\psi$, based on the maximal Smale system $T$ associated to a countable IFS with overlaps $\mathcal{S}$ as above, are given next. In Theorem 3.3 we also construct a family of measures on sub-fractals in the limit set $\Lambda$, which are related to the intricate geometry of $S$.

Theorem 3.3. In the setting of Theorem 3.2 let the countable IFS $\mathcal{S}$ with limit set $\Lambda$, and for any $s>0$ define the function $\psi_{s}: \Sigma_{I} \rightarrow \mathbb{R}$,

$$
\psi_{s}(\eta)=s \log \left|T_{i \omega}^{\prime}\left(\hat{\pi}_{2}(\eta)\right)\right|, \quad \eta \in \Sigma_{I}, \eta_{0}=i, \omega=\left.\eta\right|_{0} ^{\infty} .
$$

Then for any $s>0$,
a) $\psi_{s}$ is summable and Hölder continuous on $\Sigma_{I}$.
b) If $\mu_{s}:=\mu_{\psi_{s}}$ is the equilibrium measure of $\psi_{s}$ on $\Sigma_{I}$, and $\nu_{s}:=\pi_{0 *} \mu_{s}$ on $\Sigma_{I}^{+}$, then for $\nu_{s}$-a.e. $\omega \in \Sigma_{I}^{+}$, the measure $\nu_{s}^{\omega}:=\nu_{\psi_{s}}^{\omega}$ is exact dimensional on the sub-fractal $J_{\omega} \subset \Lambda$ and,

$$
H D\left(\nu_{s}^{\omega}\right)=\frac{h_{\mu_{s}}(\sigma)}{\left|\int_{\Sigma_{I}} \log \right| \phi_{\tau_{0} \tau_{1} \ldots \tau_{n_{\tau_{0}}(\sigma \tau)}^{\prime}}^{\prime}\left(\hat{\pi}_{2}(\tau)\right)\left|d \mu_{s}(\tau)\right|} .
$$

Proof. a) Since the maximal system $T$ associated to $\mathcal{S}$ by Definition 3.1 was shown in Theorem 3.2 to be a conformal Smale endomorphism, it follows that $\psi_{s}$ is Hölder continuous on $\Sigma_{I}$.

Also, recall the definition of $T_{i \omega}$ as a composition of maps given by the inductive relation (3.3). Also, the initial contractions $\phi_{i}: V \rightarrow V$ from $\mathcal{S}$ are defined on a compact set $V \subset \mathbb{R}^{D}$, such that $\left|\phi_{i}^{\prime}\right| \leq \alpha<1, i \in I$ on $V$. From the definition (3.3), we have the increasing sequence of integers, $n_{1}(\omega)<n_{2}(\omega)<\ldots$, and thus for any $\omega \in \Sigma_{I}^{+}$, the positive integers $n_{k}(\omega)$ satisfy the inequality,

$$
n_{k}(\omega) \geq k, k \geq 1
$$

Now from (3.3) we have $T_{i \omega}=\phi_{i \omega_{1} \ldots \omega_{n_{i}(\omega)}}$, and thus it follows from above that for any $i \in I$,

$$
\left|T_{i \omega}^{\prime}\right| \leq \alpha^{n_{i}(\omega)} \leq \alpha^{i} .
$$

Therefore, since we assumed $s>0$, it follows that,

$$
\sum_{i \in I} e^{\sup \psi_{s}[[i]} \leq \sum_{k \geq 1} \alpha^{s k}<\infty
$$

Hence the real-valued function $\psi_{s}$ is also summable in this case.
b) For any $s>0$, if $\nu_{s}:=\nu_{\psi_{s}}$ is the canonical projection of the equilibrium measure $\mu_{\psi_{s}}$ onto $\Sigma_{I}^{+}$, then in the notation of Theorem 3.2 we have for $\nu_{s}$-a.e $\omega \in \Sigma_{I}^{+}$,

$$
\nu_{s}^{\omega}:=\nu_{\psi_{s}}^{\omega}=\hat{\pi}_{2 *} \mu_{\psi_{s}}^{\omega}
$$

which is a measure supported on $J_{\omega} \subset \Lambda$. Now due to the properties of $\psi_{s}$ proved in a), one can apply Theorem 3.2 for the measure $\nu_{s}^{\omega}$ on the fiber set $J_{\omega} \subset \Lambda$, for $\nu_{s}$-a.e. $\omega \in \Sigma_{I}^{+}$. Thus we obtain the exact dimensionality of $\nu_{s}^{\omega}$ on $J_{\omega}$. Then, by using the specific expression of $T_{i \omega}$ in this case, we compute the Hausdorff dimension of $\nu_{s}^{\omega}$ by the above formula.

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